# Geothermal Mineralization I: The Mechanism of Formation of the Beowawe, Nevada Siliceous Sinter Deposit

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# GEOTHERMAL MINERALIZATION I: The Mechanism of Formation of

the Beowawe, Nevada Siliceous Sinter Deposit

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The siliceous sinter deposit at Beowawe, Nevada ABSTRACT. has a volume of  $7.71 \times 10^7$  m<sup>3</sup> and contains  $1.28 \times 10^{11}$  kg This silica precipitated from geothermal solutions silica. that welled-up along the Malpais Fault, a graben bounding Sinter formation involved three steps: (1)fault. Geothermal solutions that were saturated with quartz in the reservoir moved to the surface where they cooled and became supersaturated with amorphous silica. (2) Amorphous silica particles nucleated (heterogeneously on ferric hydroxides) to produce a colloidal suspension. (3) The amorphous silica particles agglomerated and were cemented together by amorphous silica precicitating in the embayments between the particles. At least 5.58x10<sup>14</sup> kg of geothermal solution must have been produced over time to account for the mass of silica in the Beowawe Deposit. This fluid would have carried at least  $4.95 \times 10^{17}$  kJ of heat to the surface.

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### INTRODUCTION

Siliceous sinter deposits are common features of many high temperature geothermal areas. Such areas are often characterized by boiling hot springs and geysers that produce dilute, alkaline waters. Sinter deposits occur as terraces around these features and are composed of predominately white to tan or gray, often banded, friable to dense, vitreous masses of amorphous silica. The formation of siliceous sinter requires that the geothermal solutions contain a high enough silica concentration to become saturated with amorphous silica when they cool through the 100 to 50°C range. The silica concentration of the geothermal solution is controlled by the solubility of quartz under reservoir conditions (Fournier and Rowe, 1966) so that a boiling spring depositing sinter at 100° must be fed by water containing at least 370 ppm dissolved silica. This silica concentration could only be produced in a reservoir that is hotter than  $235^{\circ}$  (fig. 1).

# 

#### Figure 1.--near here

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For a spring depositing sinter at  $50^{\circ}$ C the reservoir must be at least  $175^{\circ}$ C. Thus, siliceous sinter deposits are unlikely to form in geothermal areas where the reservoir temperature is less than  $200^{\circ}$ C.

In this paper we will describe the features of the large Beowawe, Nevada siliceous sinter deposit and discuss a

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general model for its formation. The geology of the Beowawe, Nevada area is described in detail in Struhsacker (1980) and Zoback (1979) and is sumarized here for the conveinence of the reader.

The Beowawe geothermal system, located astride the Lander-Eureka county line in the Whirlwind Valley of northcenteral Nevada (fig. 2) is a hot-water system with measured temperatures ranging from 160°C at 1830 m depth to 212°C at a depth of 2915 m (Ginn 1-13 well).

#### Figure 2.--near here

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Exposed reservoir rocks within the Beowawe area include the Paleozoic siliceous eugeosynclinal rocks of the Valmy Formation and overlying Tertiary volcanic units. The Valmy Formation consists of siliceous siltstone with significant amounts of quartzite, sandstone, bedded chert, siliceous conglomerate, and locally bedded barite. The Tertiary volcanics range in composition from basalt to dacite and 'include dikes, flows, and tuffs locally interbedded with interflow tuffaceous and alluvial horizons. In general, pervasive alteration through the Miocene volcanic section has produced a smectite-calcite-quartz-hematite assemblage near the surface that grades into a smectite-chloritecalcite-quartz-pyrite assemblage at depth (Cole, in prep.). Below the Tertiary-Ordivician contact, the rocks are weakly altered to a quartz-calcite-pyrite-chlorite-smectite-epidote

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assemblage.

Three major episodes of structural development are recognized within the geothermal area: a Paleozoic thrust faulting event, a northwest-trending graben developed during the mid-Miocene, and northeast-trending Basin and Range faulting. As a result of thrusting, allochthonous siliceous rocks of the Ordivician Valmy Formation lie in exposed fault contact upon the subjacent autochthonous carbonates along the Roberts Moutain Thrust of the upper-Devonian to lower-Mississippian Antler Orogeny (see Zoback, 1979 for schematic sections). At a much later time, major northwest-trending fractures were produced by rifting along the Oregon-Nevada These confined laterally a thick sequence of lineament. middle-Miocene, calc-alkaline to alkaline volcanic flows. The northeast-trending faults that have produced the principle topographic feature in the area, the Malpais Rim, represent the final period of Basin and Range faulting. These normal faults post-date all volcanic units. The sinter terrace has developed along the fault-controlled Malpais Rim which bounds the southeast margin of Whirlwind Valley (fig. 2).

#### THE BEOWAWE SINTER DEPOSIT

The Beowawe Sinter Deposit is located along the trace of the graben bounding fault that produced the Malpais Rim. This deposit is more than 65 m thick and 1.6 km in length and is estimated to contain more than  $1.28 \times 10^{11}$  kg of silica. Situated on this sinter terrace are numerous hot

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springs and fumaroles (fig. 2). Before intensive geothermal exploration disrupted the local hydrology, there were at least three active geysers on the sinter terrace (Nolan and Anderson, 1934 and Rinehart, 1968). Presently most of the hot water production is from uncapped exploration wells and numerous springs. The estimated combined discharge is about 400 liters per minute (Renner, et al., 1975). The temperature of the boiling springs is around 93°C which is the boiling point of water at the 5000 foot elevation of the deposit. Down hole temperature measurements and chemical geothermometers show subsurface temperatures of 200 to 250° with an average temperature of about 230°C being most The geothermal waters are dilute, probable. slightly alkaline, sodium bicarbonate-sulfate solutions (see Nolan and Anderson, 1934 and Table 1 for analyses).

Table 1.--near here

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Although the sinter terrace cannot be examined completely from top to bottom, several areas of excellent exposure do exist which can be The documented. most conspicuous feature of the terrace, in addition to its size, layering that has resulted from accretion is theof This layering can be classified into siliceous material. three categories on the basis of scale. On the scale of meters to tens of meters, the sinter has the appearance of a normal sediment, with separate units on the order of tens of

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centimeters in thickness. Some of these units are massive with fine-grained siliceous material forming a matrix that encloses larger opaline fragments or shards whose long axes lie parallel to the dip of the unit. Other units have a very fissile almost shale-like texture with an average of two to five parting planes per centimeter. This parting plane phenomena is the second class of layering. Commonly, a very dense opaline unit lies at the contact between either two massive units or a fissile unit and a massive unit. These opaline units are typically pure white but may contain black or gray bands that are only a few millimeters thick. Small solution cavities and channels are common above these units. dense opaline The third type of layering recognizable in the sinter is a millimeter scale banding. In general, each parting plane contains this finest layering but it is often obscured by white siliceous dust. These fine laminations are characteristically tinted red with iron · oxide stains in the upper portions of the sinter pile, but they become more opaline or vitreous toward the base of the sinter pile.

The attitudes of the layering are variable depending on the location on the terrace. At the top of the terrace, near the Malpais Rim, dips are relatively low, averaging less than ten degrees to the north, away from the rim. Further down and out into the pile, the units dip from 18° to 20° away from the Malpais Rim. At the bottom of the pile, the units are again nearly flat-lying with dips

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averaging less than five degrees and without a prefered orientation.

Bulk density measurements made on seven samples selected from the sinter pile range from 1.33 to 2.12 g  $cm^{-3}$ , and have an average of 1.66±0.3 g  $cm^{-3}$ . Compared to the density of silica minerals these values indicate that the void space in the sinter may be as high as 50%.

The total mass of siliceous material can be calculated if the volume of the sinter pile is known. A computer graphics technique was used to determine the area within each of the contours on the sinter deposit as shown in figure 2. It was assumed that each sub-area (12 in all) could be approximated by a rectangular block with a base at 4680 elevation. The total volume of the pile was determined by summing the volumes of each of the 12 separate blocks each having a unique area and <u>heigth</u>. This technique yielded a total volume for the pile of 7.71x10<sup>7</sup> m<sup>3</sup> (0.077 km<sup>3</sup>). If the average bulk density of the sinter is 1.66 g cm<sup>-3</sup> as calculated above, the mass of siliceous material in the pile is  $1.28x10^{11}$  kg. This is within an order of magnitude of the value of  $1.76x10^{10}$  kg that was estimated by Zoback (1979).

#### MECHANISM OF SINTER FORMATION

There are three main steps in siliceous sinter formation. First, geothermal solutions, that have equilibrated with the quartz in the surrounding rocks, move to the surface and become supersaturated with amorphous

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silica as they cool. Second, amorphous silica particles nucleate and grow to produce a colloidal suspension. Finally, the suspended amorphous silica particles agglomerate and are cemented together to form sinter.

While the geothermal fluid is in the reservoir, it is equilibrated with the quartz in the surrounding rocks (Fournier and Rowe, 1966). When some of this fluid escapes toward the surface, it travels down a steep thermal gradient so that it rapidly becomes supersaturated with quartz. However, in areas with high discharge rates, the rate of ascension of the fluid toward the surface is generally so fast that there is not time for a significant amount of silica to precipitate from the solution (Rimstidt and Barnes, 1980). At or very near the surface, the solution finally cools to the saturation temperature of amorphous silica. When the geothermal solution is discharged onto the surface amorphous silica precipitation begins.

Quartz is the thermodynamically most stable form of silica under geothermal conditions and according to simple silica-water kinetics it should always precipitate more rapidly than amorphous silica (Rimstidt and Barnes, 1980). However, in sinter deposits theleast stable silica polymorph, amorphous silica, precipitates instead. This phenomonon has often been observed in other systems and is known as Ostwald's Step Rule. Ostwald's Step Rule can be stated as: In highly supersaturated systems, the least stable (most soluble) polymorph of the solid often

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precipitates rather than the thermodynamically stable one; the metastable precipitate then recrystalizes stepwise to more and more stable polymorphs until equilibrium is achieved. While this is an empirical rule and is not followed by all systems, it seems to hold for many geochemical situations, including silica precipitation from geothermal solutions.

The phenomenon of precipitation of metastable solids is solution is so supersaturated that observed when the nucleation rather than simple molecular deposition dominates precipitation factor the process. An important in nucleation kinetics is the effect of the surface free energy of a particle of very small size on its stability. In general, the free energy of formation of a particle of radius, r, is:

 $\Delta G_{f} = -\{(4/3)_{\pi} r^{3} \Delta G_{f}^{o} / V\} + \{(4x10^{-10})_{\pi} r^{2} \sigma\}$ 

(Uhlmann and Chalmers, 1966). Here,  $\Delta G_f$  is the free energy of formation of particles of radius, r;  $\Delta G_f^o$  is the free energy of formation of the bulk solid; V is the molar volume of the solid;  $\sigma$  is the surface free energy of the solid in contact with the aqueous solution. Spherical geometry is assumed in the following discussion for simplicity. The first term is the free energy contribution of the bulk solid and the second is the free energy of formation of the interface between the particle and the aqueous solution. This relationship can be recast in terms of the solubility

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of particles of radius, r:

$$\ln (c/c^{\circ}) = \{(2/3)\sigma VB\}/\{rRT\}$$

Here, c is the solubility of particles of radius, r;  $c^{\circ}$  is the solubility of the bulk solid;  $\sigma$  is the surface free energy of the solid in contact with the aqueous solution; V is the molar volume of the solid; B is a geometric factor (16.8 for spheres). The effect of particle size on amorphous silica solubility is illustrated in figure 3.

Figure 3.--near here.

Thus, as a nuclei grows from a single molecule to a particle of macroscopic dimensions its free energy of formation first increases because its surface area is large relative to its volume, and then with further growth the volume term becomes large enough to outweigh the surface area term and the free energy of formation constantly, decreases with further growth. This means that the free energy of formation of the particle passes through a maximum where the free energy contribution from the surface is just balanced by the free energy driving the precipitation (fig. 4).

Figure 4.--near here.

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Thus, a particle with a radius larger than the critical radius will grow spontaneously. The flux of molecules (J)

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onto the growing nuclei as they pass through this process is given by:

# $J = v \exp\{B\sigma^3 V^2 / kT(RT \ln S)^2\}$

(Nielsen, 1964).  $\nu$  is a pre-exponential factor (10<sup>30</sup> has been used here); B is a geometric factor (16.8 for spherical particles);  $\sigma$  is the surface free energy of the solid in contact with the aqueous solution; V is the molar volume of the solid; k is Boltsmans constant; and S is the degree of saturation of the solution. This equation can be evaluated if the surface free energies of quartz and amorphous silica are known. Alexander (1957) has determined the surface free energy of amorphous silica to be about 45 ergs  $cm^{-2}$  by solubility measuring the of very small particles. Unfortunately, no direct measurement of the surface free energy of quartz has been made so it must be estimated using a correlation between the surface free energy of a solid and its solubility (Nielsen and Sohnel, 1971). Using 6 ppm as the solubility of quartz yields a surface free energy of about 120 ergs cm<sup>-2</sup>.

The above data were used with equation 3 to calculate the rates of homogeneous nucleation of quartz and amorphous silica as a solution containing 345 ppm dissolved silica (saturated with quartz at 230°C) cools. The surface free energies of the solids were assumed to be constant over this temperature interval. The results of this calculation are shown in figure 5.

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Figure 5.--near here.

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A flux of around  $10^{-13}$  moles kg<sup>-1</sup> sec<sup>-1</sup> is equivalent to the removal of about ten percent (35 ppm) of the dissolved silica from solution per day. Fluxes lower than this can be considered insignificant in terms of sinter formation; and in fact, considerably higher fluxes are probably responsible for most sinter precipitation. The scales at the top of the diagram show that guartz is always more supersaturated than amorphous silica. However, the rate of quartz precipitation would not become significant until the solution cooled to about 25°C. Near this temperature the rate of amorphous silica precipitation exceeds that of quartz so there is no way for significant amounts of quartz to nucleate. Thus. the only mechanism for precipitation of quartz in geothermal systems it epitaxial growth on already existing quartz grains and it has already been shown that precipitation by this mechanism is slow (Rimstidt and Barnes, 1980).

Two factors indicate that in sinter deposition amorphous silica precipitation is catalyzed by the presence of suspended solids that act as heterogeneous nuclei. First, according to the the results displayed in figure 5, if amorphous silica precipitates by a homogeneous nucleation mechanism a significant amount of silica precipitation would not occur until the geothermal solution has cooled to around  $25^{\circ}C$ . Second, careful experiments by Makrides, et al.

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heterogeneous mechanism can be estimated from equation 3 by replacing the surface free energy of the pure solid with an effective free energy for the composite nuclei. A value of 10 to 20 erg  $cm^{-2}$  would explain the scaling rates observed many geothermal power extraction systems and the in temperature range in which most siliceous sinter is formed in the Beowawe system (fig. 5). On the other hand, quartz is a crystalline material whose lattice configuration is not nearly as flexible as amorphous silica so the potential for heterogeneous nucleation of quartz is severly limited. This the near absence of primary quartz explains in most geothermal scale and sinter deposits. Note that when the effective surface free energy is very low the contours of nucleation rate as a function of temperature (fig. 5) are very steep. This means that once the solution cools to the temperature of critical saturation nearly all of the silica nucleates over a temperature interval of less than 5°. Thus, the onset silica deposition is abrupt and is confined to a narrow zone around the vent where the water rapidly cool's through this critical temperature interval.

The most important substances that catalyze amorphous silica nucleation are ferric hydroxides. This is born out by observations of the rate of formation of siliceous scales in geothermal energy extraction systems. Experience with silica scaling in the geothermal power extraction systems of the Imperial Valley of California shows that the severity of scaling increases toward the north where both the geothermal

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solutions and the scale become more iron rich. Furthermore, scale formation experiments by Rothbaum et al. (1979) using geothermal fluids from Broadlands and Wairakei showed that aeration of the geothermal fluids produced a significant increase in the rate of scale formation as a result of the oxidation of ferrous iron to ferric iron. The ferric iron then precipitates from the neutral to slightly alkaline solutions as amorphous ferric hydroxides. Because ferric hydroxides cannot form in the unaerated solutions that are in the vents and guiet pools, these features do not become filled and choked by sinter deposition, but rather, most of the sinter deposits occur on the flanks of these features where the flowing water is both aerated, to produce ferric hydroxide nuclei, and cooled rapidly through the critical temperature range.

The final step in the formation of siliceous sinter is the agglomeration and attachment of the suspended particles to the surface of the sinter deposit. This process is effected by the direct precipitation of dissolved silica into the contact zone between two silica particles. Silica precipitation is enhanced in this zone because silica solubility is considerably reduced in areas with very small negative radii of curvature (Iler, 1973)(fig. 3). Thus, precipitating silica rapidly fills the embayment between any silica particles that come into contact and cements them together.

This cementation model explains the sediment-like

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layering and the various textures of the sinter deposit. Under high flow rates or turbulance, such as near the vent, most suspended particles are swept off of the sinter surface before they can adhere. However, the few particles that do stick to the sinter become heavily overgrown by silica precititating directly from solution to produce a deposit of dense, vitreous sinter. Thus, the sinter near the springs and geysers is hard and dense. On the other hand, under low water velocities many more suspended particles can accumulate on the sinter surface and each layer of particles is covered by a successive one before they have been thoroughly cemented together. Thus, a more porous, chalky sinter is formed at a distance from the vents.

# TOTAL DISCHARGE OF THE BEOWAWE SYSTEM

The total discharge of the Beowawe system over time can be estimated from the mass of the sinter deposit and the amount of silica deposited per kilogram of solution (P, see fig. 6) using the following equation:

#### D = M/P

Figure 6.--near here.

The bulk density of the sinter is estimated to be 1.66 g  $cm^{-2}$  and the total volume of the sinter deposit is estimated as  $7.7 \times 10^{13}$  cm<sup>3</sup>; thus, the mass of the sinter deposit is about  $1.28 \times 10^{14}$  g. P is the difference between the

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solubility of quartz at the reservoir temperature and that of amorphous silica at 25°C. If the reservoir temperature were 230°C throughout the life of the system, about 5.58x10<sup>14</sup> kg of water would have been required to produce the Beowawe sinter deposit. It is possible that in the past the reservoir temperature of the Beowawe system was higher so the discharging water would have produced more sinter per unit mass. Assuming a reservoir temperature of 300°C as a reasonable maximum, the minimum amount of water necessary to produce this sinter deposit is around 2.66x10<sup>14</sup> kg.

The size of the sinter deposit can also be used to estimate the total amount of heat carried to the surface by the geothermal waters. This can be done because both the amount of heat and the amount of silica brought to the surface per unit mass of water are a function of the reservoir temperature. Thus, the total amount of heat transported to the surface (H) over the life of the system can be found from:

# $H = M(\Lambda H/P)$

The enthalpy difference ( $\Delta H$ ) between pure water at saturation pressure and the temperature of the reservoir and the enthalpy of water at 25°C was calculated from Keenan, et al. (1978). The ratio of  $\Delta H/P$  has been calculated for a range of temperatures and is displayed in figure 6. If the reservoir temperature were 230°C the total heat production was  $4.95 \times 10^{17}$  kJ. On the other hand, if the reservoir

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temperature were 300°C the heat production was  $3.30 \times 10^{17}$  kJ. Thus, the total amount of heat brought to the surface by the geothermal fluid that produced the Beowawe sinter deposit is around  $4 \times 10^{17}$  kJ. Around 1.2 kJ g<sup>-1</sup>, of heat is released by cooling a basaltic intrusion from its liquidus temperature to  $250^{\circ}$ C (Norton and Cathles, 1979). Based on this, the minimum size of a basaltic intrusion that could be responsible for furnishing the heat to produce the Beowawe sinter deposit is  $3.3 \times 10^{17}$  g (115 km<sup>3</sup>).

#### CONCLUSIONS

The sinter mound at Beowawe has many features similar to those found in sedimentary rocks. First, it shows bedding on a scale of tens of centimeters that is similar to sedimentary bedding. On a smaller scale it shows partings similar to those in shales. Finally, on a millimeter scale are fine bands. When the orientation of these features is observed in map scale a delta pattern emerges. It is likely that the sinter deposition was initiated at numerous points along the base of the Malpais Fault scarp and that the resulting sinter mounds coalesced to form a continuous deposit. These nearly flat-lying beds (dips less than 10°) are analogous to the topset beds of a delta. After these original outlets became choked with sinter, new conduits were formed by water flowing laterally through the sinter pile to discharge along its flanks. This activty on the flanks of the deposit produced the most steeply dipping (10-20°) layering that is dipping away from the source area

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similar to foreset beds. Although geothermal sources may have occured further out in the valley their contribution to the sinter mass was apparently negligable compared with the amount of sinter deposited from the fluids produced through the Malpais fault trace. The beds in this area show low dips with no particular orientation and may be compared to bottomset beds. At the present time, fumarole and hot spring activity is almost exclusively restricted to the top and southwest end of the terrace.

The mechanism of formation of the siliceous sinter in the Beowawe deposit involves three steps. First, geothermal solutions that are saturated with quartz in the reservoir move to the surface and cool. When the solutions have cooled to a critical temperature amorphous silica nucleates. This nucleation is catalyzed by other suspended solids in geothermal water, particularly amorphous the ferric hydroxides. Finally, the particles of amorphous silica are agglomerated and attached to the sinter surface. This is a result of the rapid precipitation of silica around the contact point between two particles so as to minimize the surface free energy of the system. This mechanism is probably also responsible for the formation of siliceous scale in geothermal power systems.

The size and bulk density of the Beowawe deposit can be used to estimate that the amount of water discharged by the system was around  $5 \times 10^{14}$  kg. Furthermore, this water carried about  $4 \times 10^{17}$  kJ of heat to the surface. Although

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the actual heat source is unknown, if the heat were furnished by a cooling basaltic intrusion, it would have to contain at least 115 km<sup>3</sup> of basalt to supply enough heat to account for the size of the Beowawe sinter deposit.

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The	chemical	comp	posit	tion	of	fluids	and
•	sinter	from	the	Beov	vawe	e area	

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	(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)
	Hot Sp.	Hot Sp.	Hot Sp.	Well	Well	Sinter	Sinter	Sinter
T, °C	89	89	95	132	160	<u> </u>		
Τ, ∘C(9)	192	186	203	217	240			
pH (10)	8.65	8.45	8,55	9.3	9.17			
Na(ppm)	219	206	204	214	277	2841	4948	3672
K (ppm)	18	14	24	9	35	1229	7297	2059
Ca(ppm)	12	1	1	13	2.5	2258	2287	2516
Mg(ppm)	1	<0.5	<0.5		0.3	289	772	326
Fe(ppm)	0.66	<0.025	<0.025		0.53	668	11193	1337
Al(ppm)	<0.625	<0.625	<0.625	0.2	1.1	1276	29059	3287
Ti(ppm)	<0.125	<0.125	<0.125			54	1906	138
P (ppm)	<0.625	<0.625	<0.625			22	271	26
Sr(ppm)	0.03	0.04	0.03			82	181	108
Ba(ppm)	0.63	<0.625	<0.625			27	645	54
Mn(ppm)	<0.25	<0.25	<0.25			23	85	31
W (ppm)	0.2	0.13	<0.125			<1200	<1200	<1200
As(ppm)	<0.625	<0.625	<0.625		<0.625	<25	<25	<25
Hg(ppb)						50	4395	137
Li(ppm)	1.52	1.43	1.50		1.94	13	34	13
Be(ppm)	<0.005	<0.005	<0.005			3.6	3.4	5.4
B (ppm)	2.0	1.8	1.9	1.0	2.0	<400	<400	<400
SiO <sub>2</sub> (ppm)	236	218	274	329	436			
HCO3(ppm)	345	340	159	33	175	·	_	
CO <sub>3</sub> (ppm)			196	246	92			
S0 <sub>4</sub> (ppm)	99	92	107	89	76			
Cl (ppm)	31	19	36	50	67			
F (ppm)	18	22	15 ·	6	12.2	·		
TDS(ppm)	958	898	1000	989	1240			

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Notes: (Table 1)

- (1) Hot spring sampled 12-3-81, location 1 in Figure 2.
- (2) Hot spring sampled 12-3-81, location 2 in Figure 2.
- (3) Boiling hot spring sampled 12-3-81, location 3 in Figure 2.
- (4) Vulcan well 2, analyses given by Cosner and Apps (1978), see Figure 2 for map location.
- (5) Chevron well 85-18, analyses given by Iovenitti (1981(, see Figure 2 for map location, well depth is 1804m (5920 feet)
- (6) Siliceous sinter sampled 12-3-81, location A in Figure 2.
- (7) Siliceous sinter (massive) sampled 12-3-81, location B in Figure2.
- (8) Siliceous sinter (fissil) sampled 12-3-81 location B in Figure 2.
- (9) Temperature calculated from the silica geothermometer (Fournier and Rowe, 1966).
- (10) pH of the hot springs measured at temperature of spring, for wells, pH measured at ~20°C.

- Fig. 1. Quartz and amorphous silica solubility as a function of temperature. A geothermal solution in equilibrium with quartz at 238°C will become supersaturated with amorphous silica when it is cooled to less than 100°C.
- Fig. 2. Map of the Beowawe sinter deposit.
- Fig. 3. Amorphous silica solubility at 25°C as a function of the radius of curvature of the particles. For particles with a positive radius of less than 0.05 micrometers the solubility is noticably greater than the bulk solubility. If the negative radius of curvature is less than 0.05 micrometers in the embayment between two particles, the solubility is <u>less</u> than the bulk solubility and the rate of silica precipitation will be accelerated in this zone. This causes the particles to be cemented together. (Diagram modified from Iler, 1973)
- Fig. 4. The free energy of formation of a nucleus as it grows from a single molecule through the critical size (r\*) to become a stable particle. Note that the free energy barrier is smaller in the case of heterogeneous nucleation because the effective surface free energy term in equation 1 is smaller.
- Fig. 5. The flux of aqueous silica (J) (both polymers and monomers) onto solid silica particles forming in the solution as a function of temperature for a solution containing 345 ppm dissolved silica. The the left illustrate homogeneous curves on while those on the right are for nucleation heterogeneous nucleation of amorphous silica where the effective surface free energy is 10 and 20 ergs cm<sup>-2</sup>.
- The amount of amorphous silica (P) that can Fig. 6. precipitate from a geothermal solution cooled to 25°C as a function of the reservoir temperature. The solubility of amorphous silica at 25°C is 116 ppm and the solubility of quartz was calculated from the temperature functions reported in Rimstidt and Barnes (1980). Also, the amount of heat brought to the surface by the geothermal solution per gram of silica that it deposits (AH/P) as a function of the reservoir temperature. The enthalpy difference between the water in the reservoir at saturation pressure and 25°C was calculated from the values given in Keenan, et al. (1978).





Fig 2



Fig 3



# **FREE ENERGY**

tig 4



Fig 5



FIG D